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**Ultrasonic Physical Modeling of Seismic Wave
Propagation from Graben-Like Structures**

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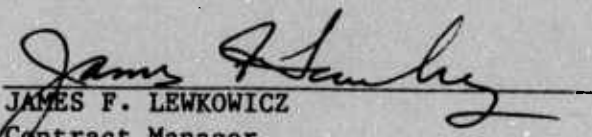
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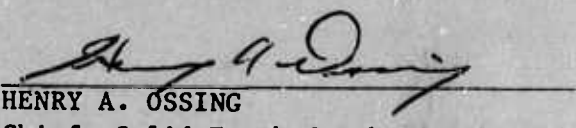
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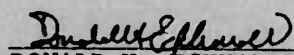
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cy part (200-400 KHz) of the signal. If we scale the results so that they correspond to a cylindrical basin roughly the size of Yucca Flat, this means that frequencies analogous to 3-10 s in the Earth appear to be amplified relative to lower frequencies. When the source is excited in the graben in an off-center position, a radiation pattern is established, with amplitude varying by a factor of 2 or more. Material effects appear to be accentuated when the source is excited off-center.

Experiments on the scale model of Yucca Flat yield, in general, less dramatic results. The loss of symmetry apparently leads to less opportunity for focussing, and amplifications are not as great. Waveforms do not vary as much azimuthally in character or shape as they do in the cylindrical case, nor do they vary as much with source position within the graben. The most dramatically different waveforms are obtained when the source is excited over the deepest portion, in the southern part of the structure. This is outside the area corresponding to the portion of the basin where most explosions have been detonated. The effect is also enhanced for wave propagation directions to the northwest or southeast. Although effects are not as dramatic as they are in the case of the cylindrical graben, there is still a relative amplification of what corresponds to waves of period 2 s or shorter in the Earth, caused by the presence of the structure.

In the real Earth, Yucca Flat is not embedded in a homogeneous half space, but in a more complex multilayered structure, this possibly having a significant effect on seismic waveforms. It would be beneficial to conduct experiments involving grabens embedded in multilayered structures, and themselves filled with a multilayered material.

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1. Introduction

The overall purpose of the Ultrasonic Physical Modeling Program at the Rockwell International Science Center is to model seismic wave propagation in the Earth using ultrasonic wave propagation in scale laboratory models. By using well-calibrated sources and receivers, our hope is to shed light on the effects of complex structure and geology on the propagation of seismic waves, and thus aid the national research effort in seismic monitoring of nuclear explosions. The intent is to complement numerical modeling, providing insight and guidance in complex situations where such modeling may not yet be feasible, owing to limitations in computer power.

In this report, we address the general problem of a nuclear explosion source region which has material properties significantly different from those of the surrounding seismic wave propagation medium. Such a situation exists, for example, in the case of explosions set off in Yucca Flat at the Nevada Test Site. The existence of a source region with differing material properties from the surrounding medium can have considerable effects on recorded surface wave amplitudes, as has been shown by some numerical studies (e.g., Regan and Glover, 1985). This in turn has implications for yield estimation, and possibly for discrimination.

2. The Receiver

It is absolutely imperative in a study of this kind to have a receiver with a well-known response. We use an NBS-type conical transducer (Proctor 1980, 1982a,b) manufactured by Industrial Quality, Inc.; it is shown in Fig. 1. This transducer is a vertical component displacement sensor with a 1 mm contact area, and a very flat response. The element is piezoceramic, and it is coupled to a large brass backing which effectively eliminates resonances, as well as minimizing coherent reflections back into the element. Figure 2 shows typical response curves for this type of transducer, sent to us by NBS. The response is flat enough that when we look at a signal from this transducer, we can consider that we are looking essentially at raw vertical component displacement.

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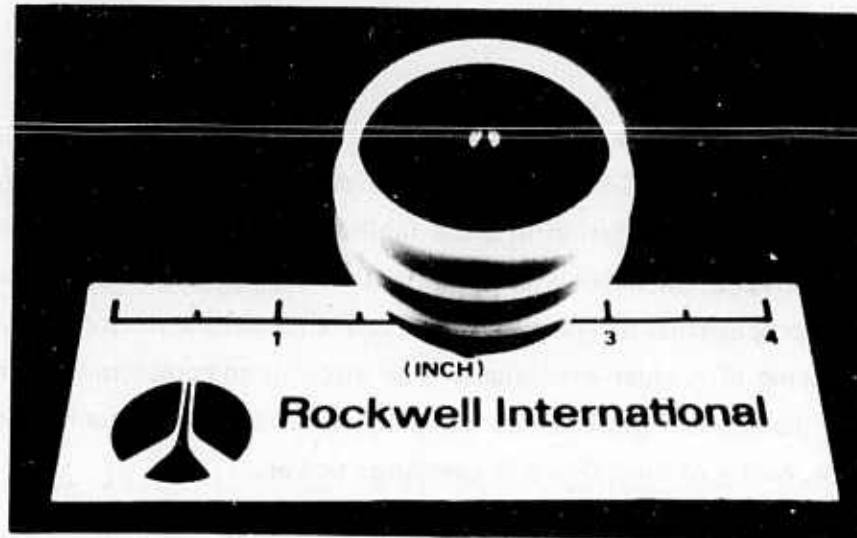


Fig. 1 The NBS-type conical transducer used in this study, showing the point-like probe.

3. The Source

Just as important as having a well-characterized receiver is having a well characterized source. The source we use is a simple one, but it is quite effective. Basically, we achieve a step-function point unloading of the surface by breaking a pencil lead on it. This is a variant of the well-known breaking-glass-capillary source used by the NBS, and is discussed in detail by Hsu and Hardy (1978). Figure 3 shows a picture of the source assembly, and Fig. 4 shows the source time function of the breaking pencil lead, obtained via deconvolution by Hsu and Hardy. The apparent noisiness in the response is due to the deconvolution process. The source approximates a step function; actually it is a ramp, but the rise time of the ramp is less than $1 \mu\text{s}$.

4. Lamb's Problem

Figure 5 shows the result of a measurement made by setting off the source on the gabbro "halfspace", and recording the signal received by the transducer 200 mm away

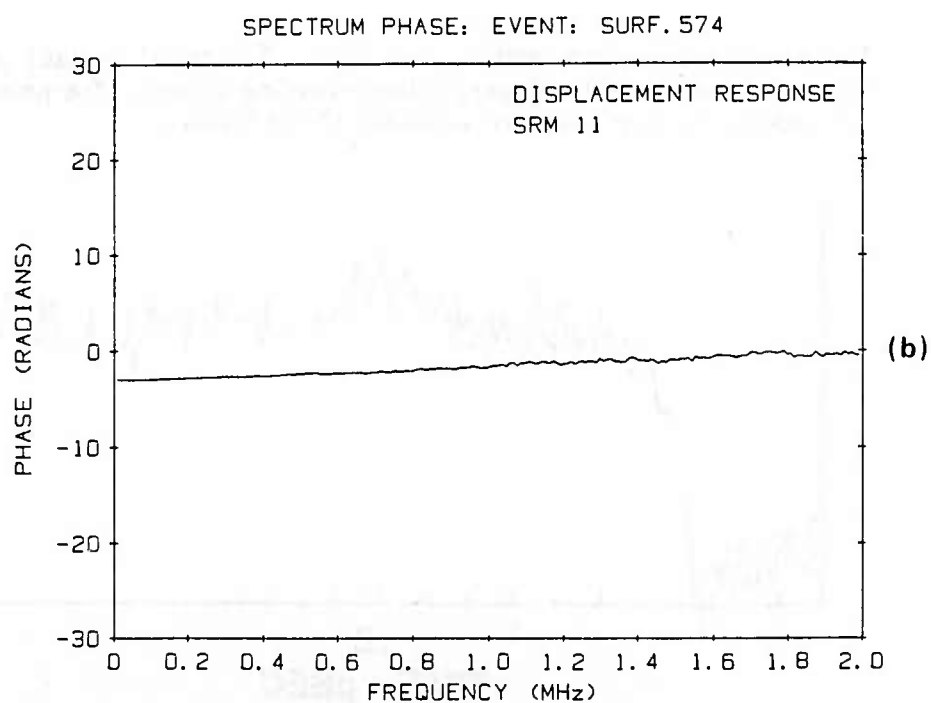
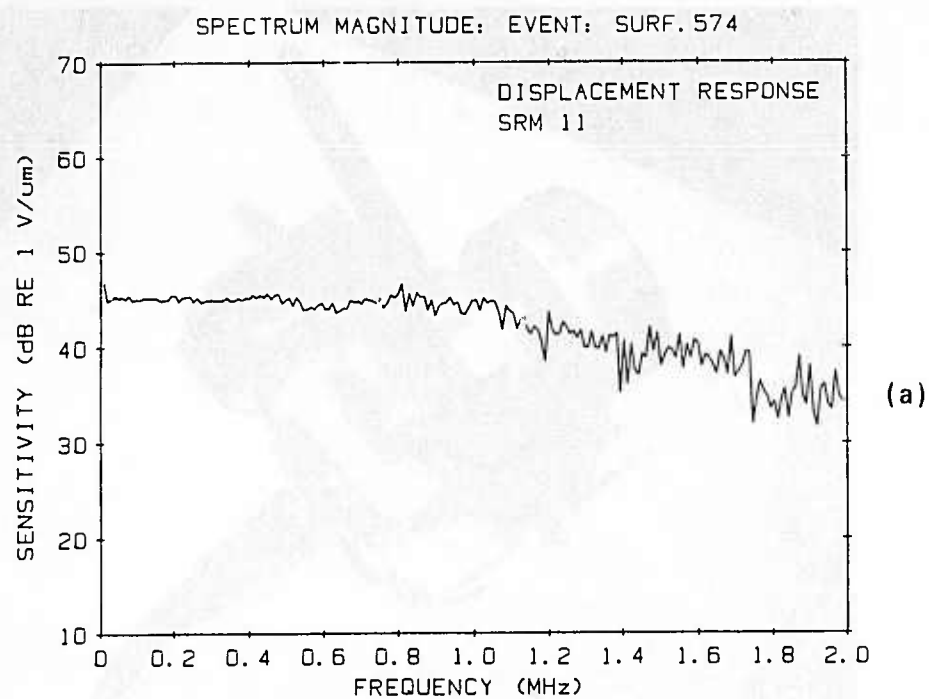


Fig. 2 Typical displacement response curves for the NBS-type conical transducer.
a) Amplitude. (b) Phase. The receiver is close to a true displacement sensor.

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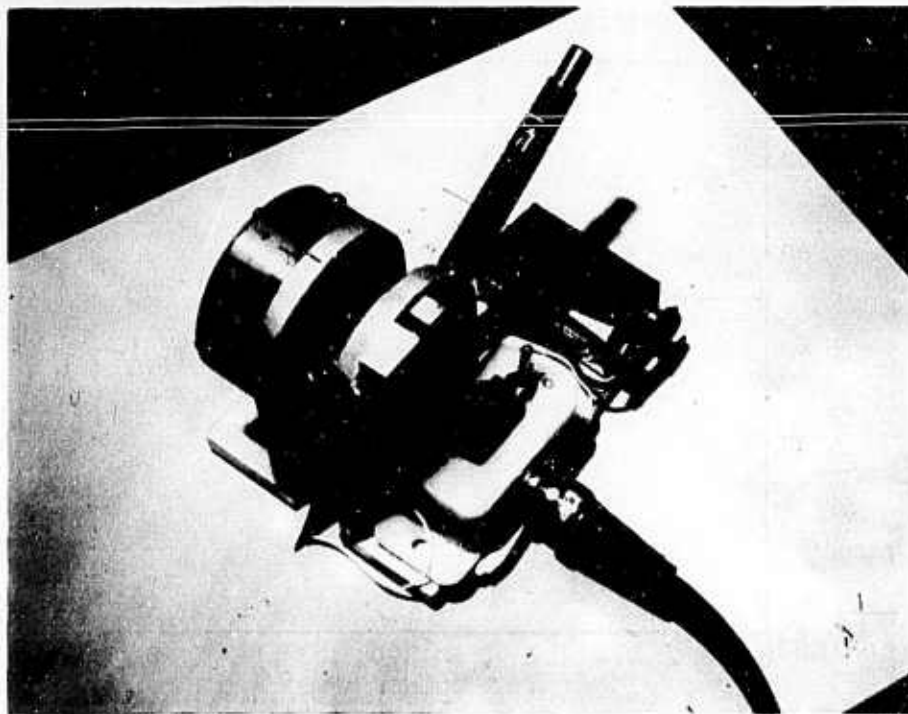


Fig. 3 The pencil-lead source used in this study. Electrical contact is broken when the pencil lead breaks, triggering the recording system. The pencil-lead source corresponds to step function unloading of the surface.

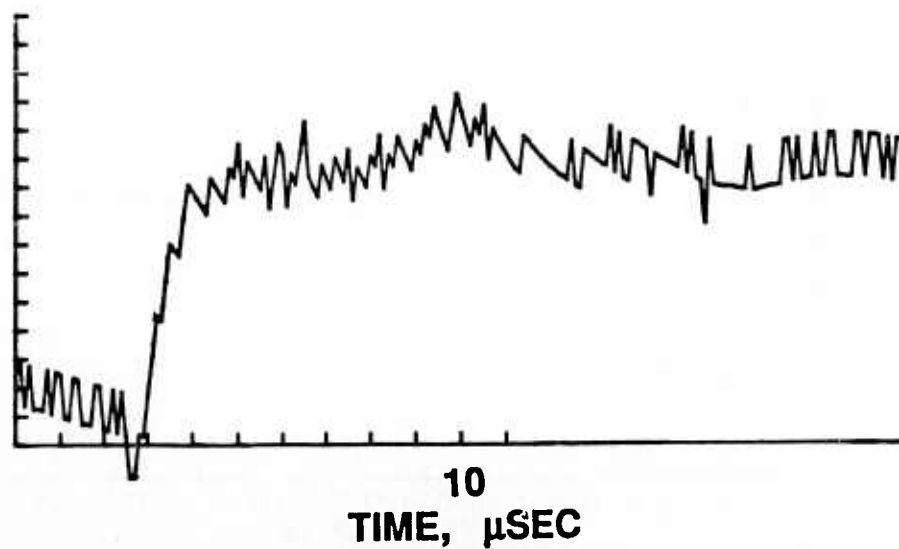


Fig. 4 Source-time function of a pencil lead source, obtained by Hsu and Hardy (1978) by deconvolution. Some spurious structure has been introduced by the deconvolution.

(the standard distance for all the measurements presented in this report). The displacement record is essentially a solution of the classic Lamb's problem (see e.g., Miklowitz, 1978; Mooney, 1974; Breckenridge et al, 1975) for a point force on a surface. Figure 5 shows, for comparison, the result of Boler et al, (1984) for a similar setup, using a breaking-glass-capillary source and a true displacement transducer. The results are very similar in appearance to ours. Boler et al include the theoretical response computed from Lamb's solution. The first arrival P wave is very small in the theoretical solution, and is very small in Boler et al's measurements. In our results, there is only a hint of it, as a minor inflection before the onset of the large signal. The large signal observed in both our record and in Boler et al's is, of course, the S wave followed by the Rayleigh wave.

5. The Cylindrical Graben Model

As a first step toward studying this problem, we have studied a cylindrical low velocity "graben," or plug, embedded in a high velocity medium (Fig. 6). The high velocity medium is a fine-grained gabbro with $V_p = 6.2$ km/s, $V_s = 3.6$ km/s, and $V_R = 3.3$ km/s. The plug is filled with lower velocity materials, whose properties are shown in Table 1.

Table 1
Properties of Modeling Materials

Material	Longitudinal Velocity V_p , km/s	Shear Velocity V_s , km/s	Rayleigh Velocity V_R , km/s	Poisson's Ratio ν	Density ρ , g/cc
Crystalbond 504 (Aremco Prods. Inc.)	2.407	1.096	1.01	0.369	1.32
HPAL3	3.287	1.742	1.61	0.305	2.01
Gabbro	6.200	3.623	3.33	0.240	2.97

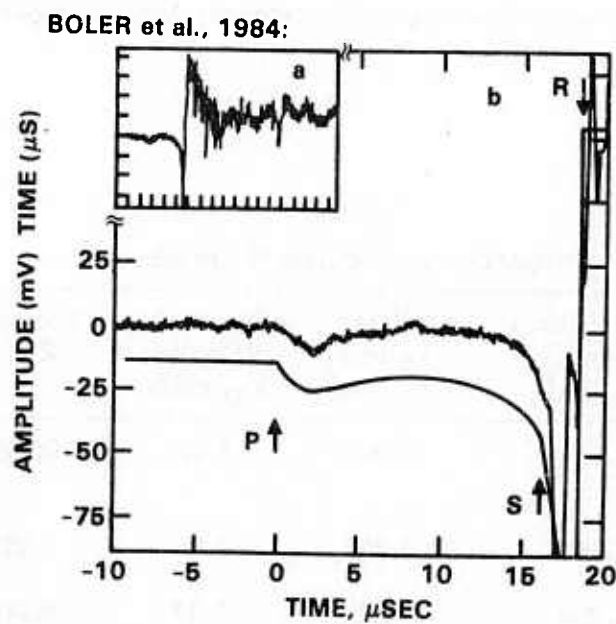
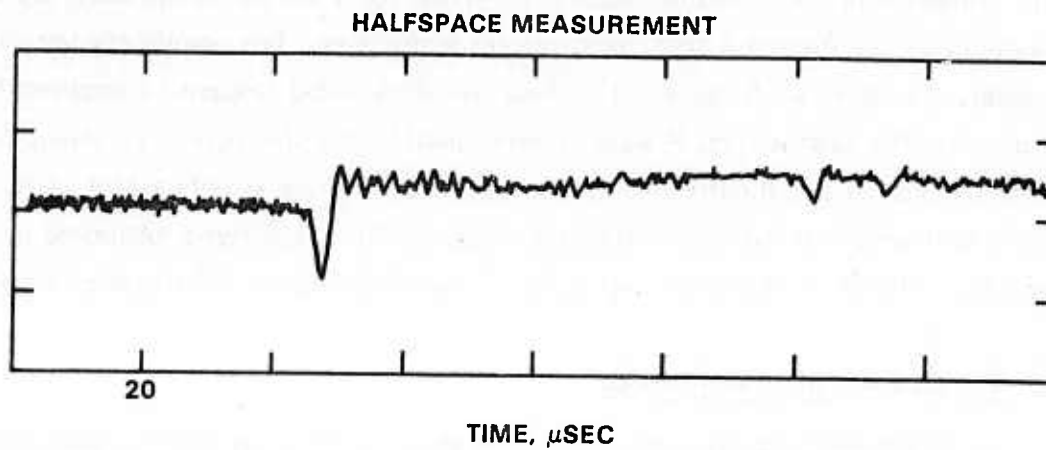


Fig. 5 Signal observed by actuating the source on the gabbro "halfspace." Similar signals obtained by Boler et al (1984) are shown for comparison.

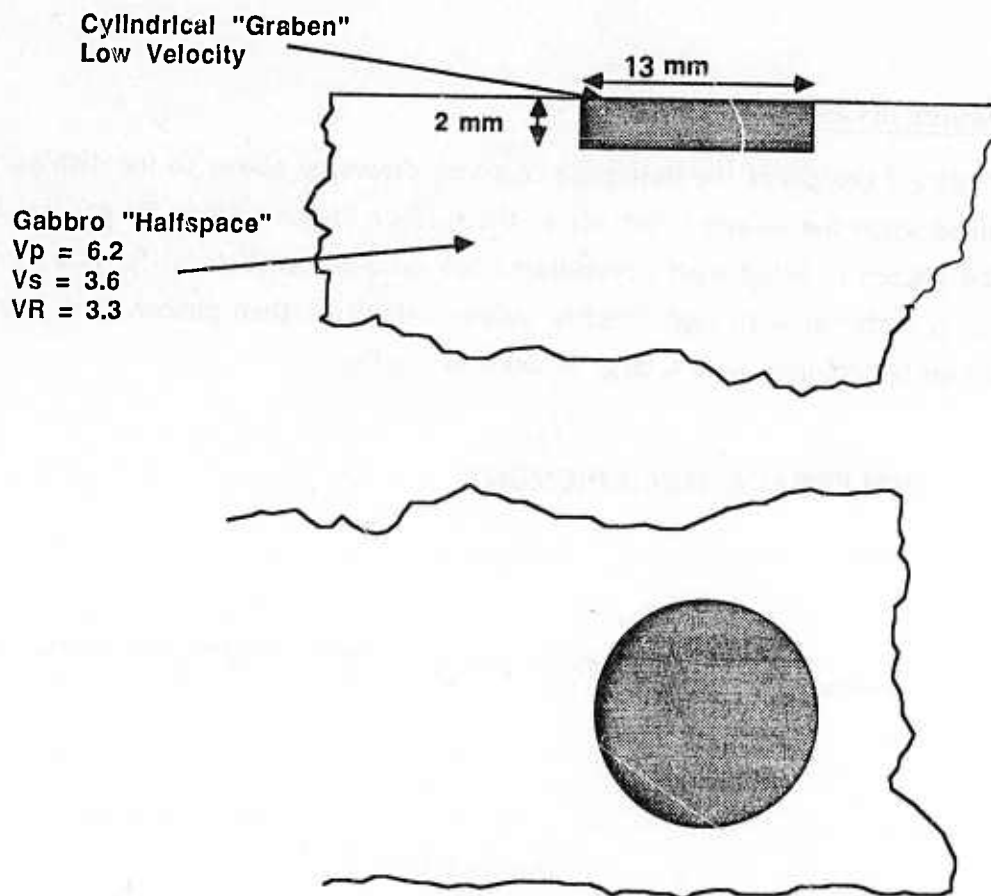


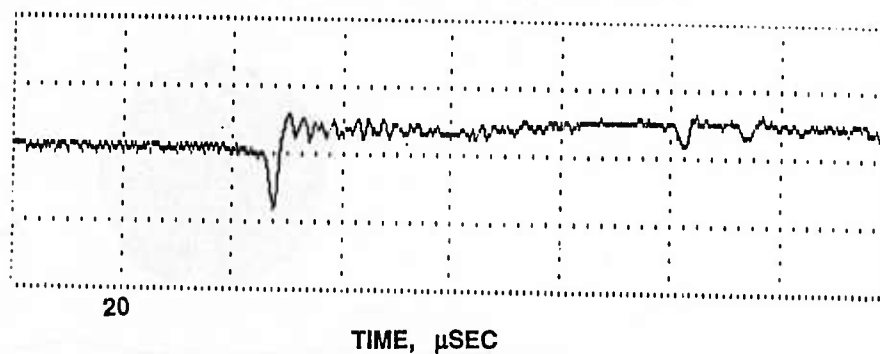
Fig. 6 The model of a cylindrical graben filled with low velocity material, embedded in a fine-grained gabbro "halfspace."

It is important to have a good idea of the scale factors involved. Taking Yucca Flat as a rough guideline, we may say that a graben of interest in the Earth is roughly $L^e = 20$ km in diameter. If the source material in the Earth has a Rayleigh wave velocity $V_R^e = 1.2$ km/s, then a 20 s Rayleigh wave in the Earth has a wavelength $\lambda_R^e = 24$ km $\approx L^e$. Now, the model graben has a diameter L^m of 13 mm. We would like to know the frequency in the model of the Rayleigh wave analogous to a 20 s Rayleigh wave in the Earth. The wavelength of this analogous wave in the model graben must be roughly equal to the graben diameter, i.e., $\lambda_R^m \approx L^m$. Since V_R^m ranges from roughly 1 to 1.6 km/s, this means that the frequency ranges from roughly 80 to 120 kHz, depending on the material in the graben. Hence, Rayleigh waves of 80 to 120 kHz in the model are analogous to 20 s Rayleigh waves in the Earth.

5.1 Source in Graben, Centered

Figure 7 compares the halfspace response discussed above to the displacement signal obtained when the source is set off at the surface of the cylindrical graben, in the center. The graben is filled with Crystalbond 504 (also referred to as "crystal wax" in the figures), a material with significantly slower velocities than gabbro (see Table 1). The signal is quite complex, with a large amount of ringing.

HALFSPACE MEASUREMENT



SOURCE AT GRABEN CENTER GRABEN FILL = CRYSTAL WAX

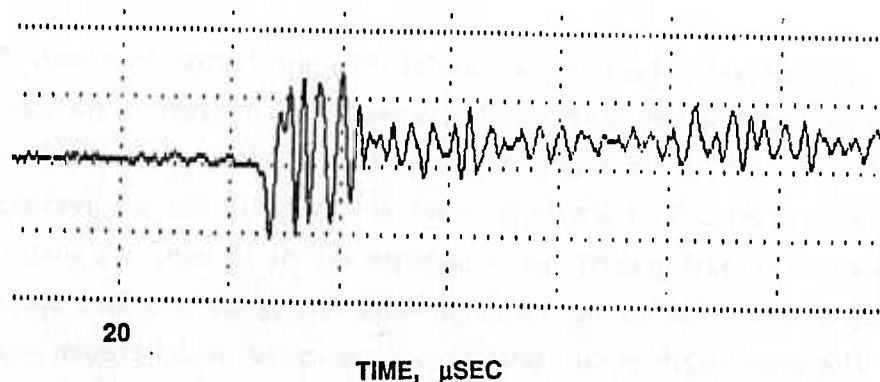


Fig. 7 Comparing the halfspace signal (also shown in Fig. 5), with the signal observed when the source is actuated at the center of the surface of the cylindrical plug ("graben") filled with Crystalbond 504. Source-receiver distance is 200 mm. Vertical scale is 100 mV per division.

Energy which, when the source is set off on the halfspace, goes downward and is not recorded at the surface, is now trapped and redirected by the graben structure.

Figure 8 compares the results from a centered source in the graben for two different fill materials. The top trace is a copy of the signal discussed immediately above, where the graben is filled with Crystalbond 504. The bottom trace is for a graben filled with HPAL3, an aluminum-filled resin with faster velocities than Crystalbond 504, but slower velocities than gabbro. As might be expected, the amplitude of the ringing is smaller than in the case of Crystalbond 504. As the material property contrast increases between the graben and the surrounding medium, the observable effects of ringing appear to increase.

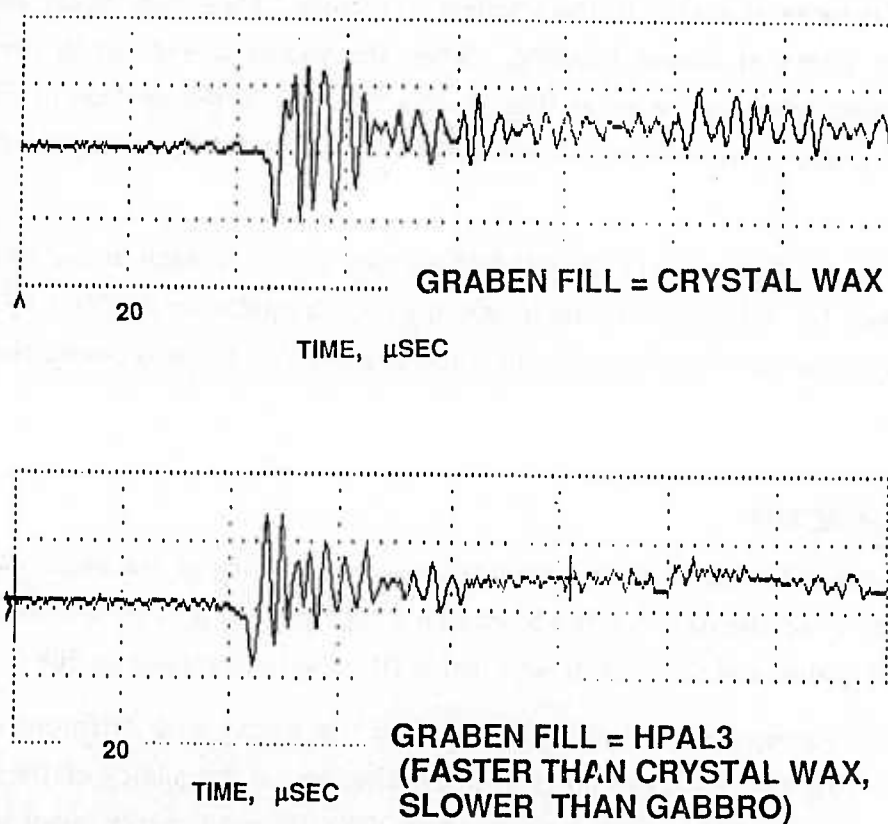


Fig. 8 Comparing signals from a source at the graben center. Top trace is for a graben filled with crystal wax (same as Fig. 7); bottom trace is for a graben filled with HPAL3. HPAL3 is faster than Crystalbond ("crystal wax"), but slower than gabbro. Source-receiver distance is 200 mm. Vertical scale is 100 mV per division.

5.2 Source in Graben, Off-Center

Figure 9 shows signals obtained when the source is actuated in the graben in various off-center positions. The relative position of source and receiver is shown schematically in plan view beside each trace. In each case, the source is actuated along a diameter, halfway between the center and the rim of the graben. (It is easy to see that this is as if the source were kept in one of the three positions, and the receiver were moved around.) Clearly, an off-center source produces a radiation pattern. Both the shape and amplitude of the signal depend on the relative position of source and receiver. The trace with the largest amplitude has a maximum peak-to-peak amplitude about twice as large as that with the smallest amplitude. These results are fairly easy to rationalize in terms of simple focusing. When the source is excited in the off-center position furthest from the receiver (Fig. 9, top trace), a larger portion of the boundary between the graben and the rest of the medium is illuminated in the direction of the receiver.

Figures 10 through 12 show the off-center signals in each of the three positions just discussed, for different fill materials (again, Crystalbond 504 and HPAL3). The effect on amplitude of the different fill materials appears to be accentuated in the off-center cases.

5.3 Voiceprints

Figures 13 and 14 show an interesting presentation of the data. What is shown is a "voiceprint" of the data for the source on a halfspace (Fig. 13), and the data for the source in the graben center when the graben is filled with Crystalbond 504 (Fig. 14).

The voiceprint is obtained by filtering the traces with different bandpass filters, and plotting the results in order of increasing center frequency of the bandpass filter. In this case, the filters have a passband of 200 kHz, and the increment in center frequency between traces is 40 kHz. Thus, the bottom trace shows the data filtered from 0-200 kHz (center frequency 100 kHz), the next trace up shows the data filtered from 40-240 kHz (center frequency 140 kHz), the next trace after that shows the data filtered from 80-280 kHz (center frequency 180 kHz), etc. What results is essentially a frequency-time plot. (Note that the traces are also rectified and low pass filtered, to avoid spurious wiggles resulting from the increasing center frequency of the bandpass filter.)

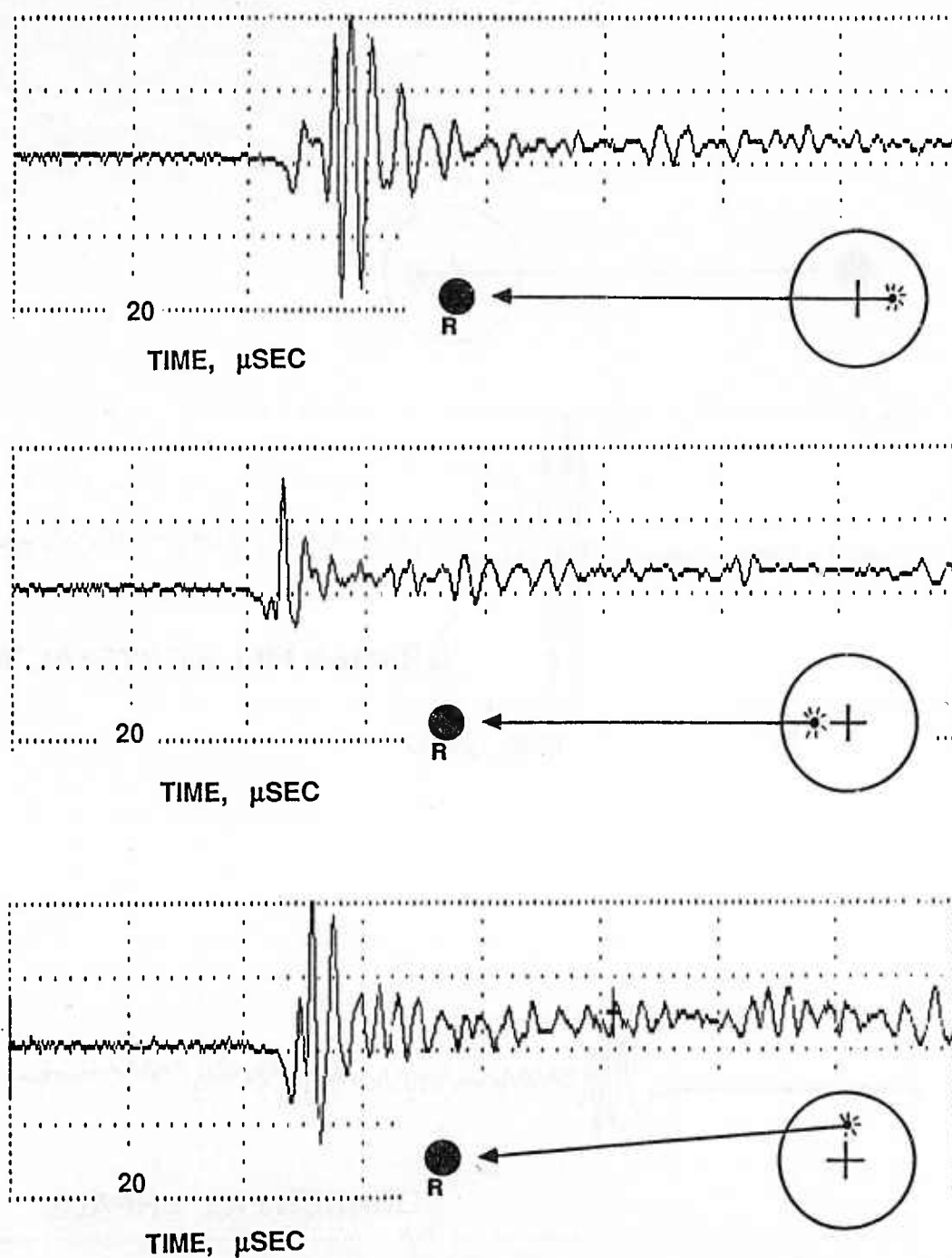


Fig. 9 Signals from sources actuated off-center in a graben filled with Crystalbond 504 ("crystal wax"). Each trace is accompanied by a plan view showing the relative positions of source and receiver. Distance from graben center to receiver is 200 mm. Vertical scale is 200 mV per division for the two top traces, and 100 mV per division for the bottom trace.

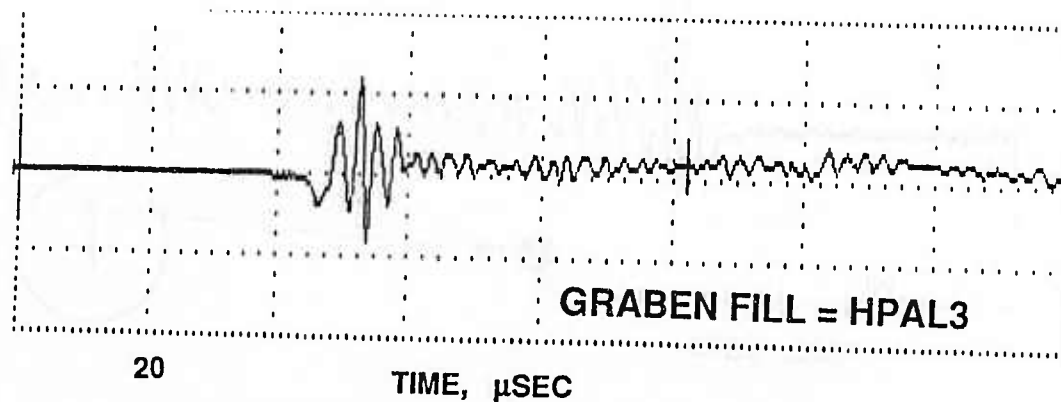
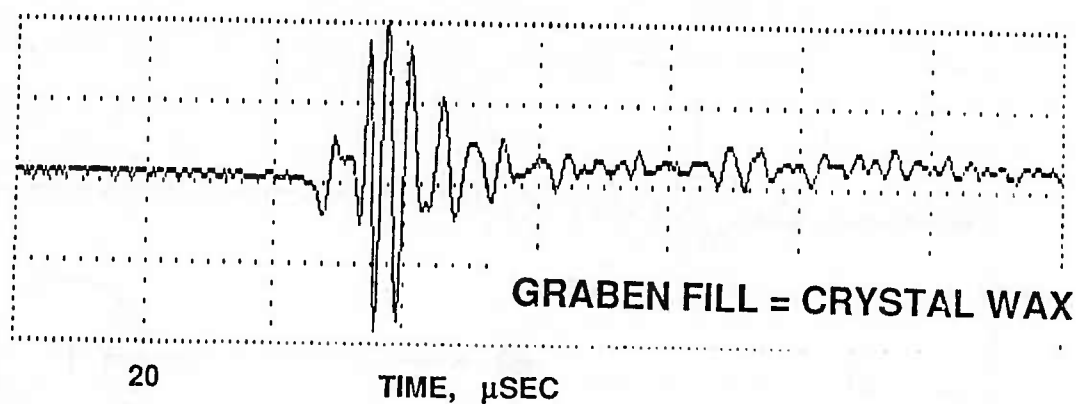


Fig. 10 Signals for one of the off-center positions in Fig. 9, for a graben filled with Crystalbond 504 ("crystal wax") and a graben filled with HPAL3. Distance from graben center to receiver is 200 mm. Vertical scale is 200 mV per division.

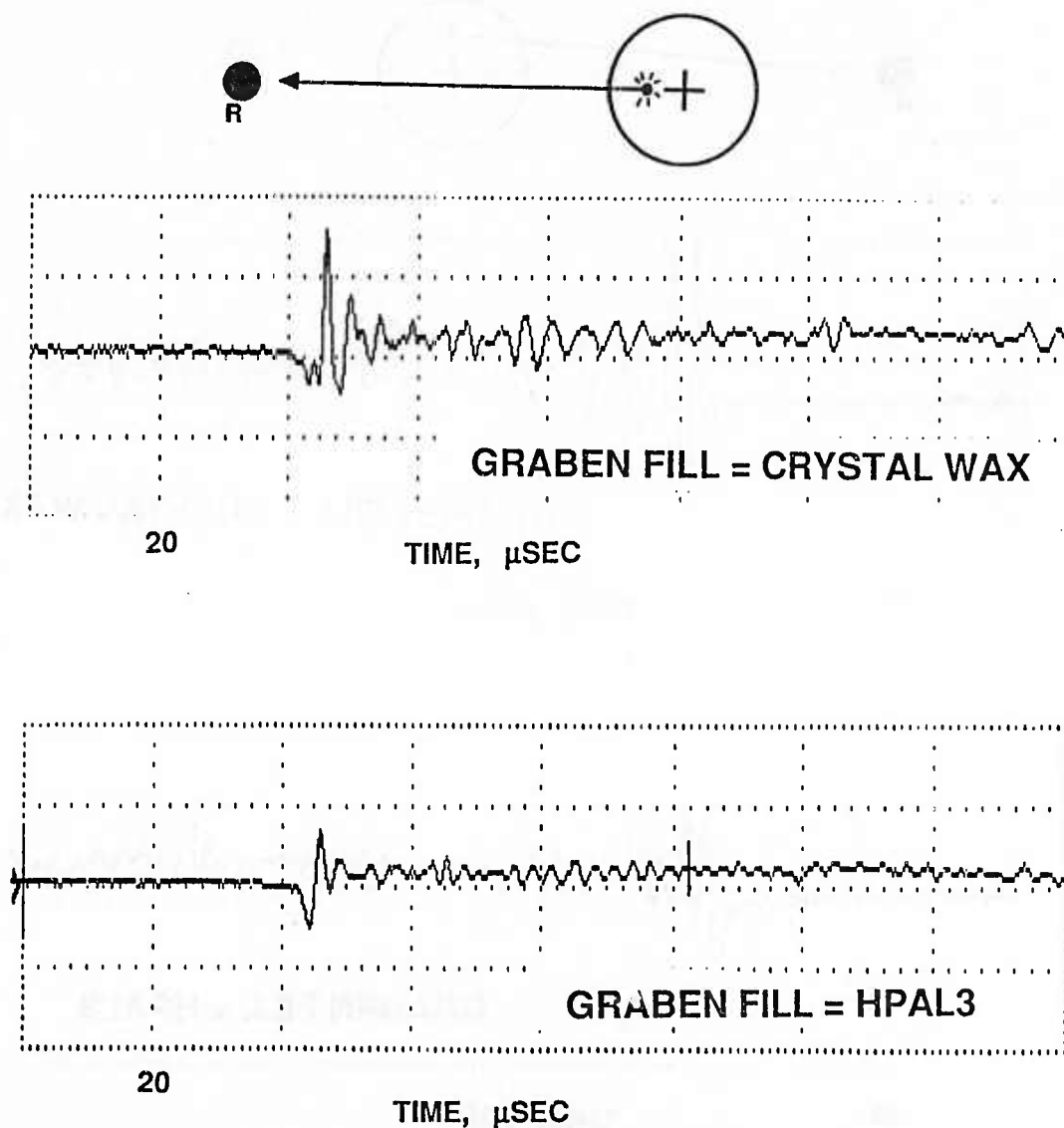


Fig. 11 Signals for one of the off-center positions in Fig. 9, for a graben filled with Crystalbond 504 ("crystal wax") and a graben filled with HPAL3. Distance from graben center to receiver is 200 mm. Vertical scale is 200 mV per division.

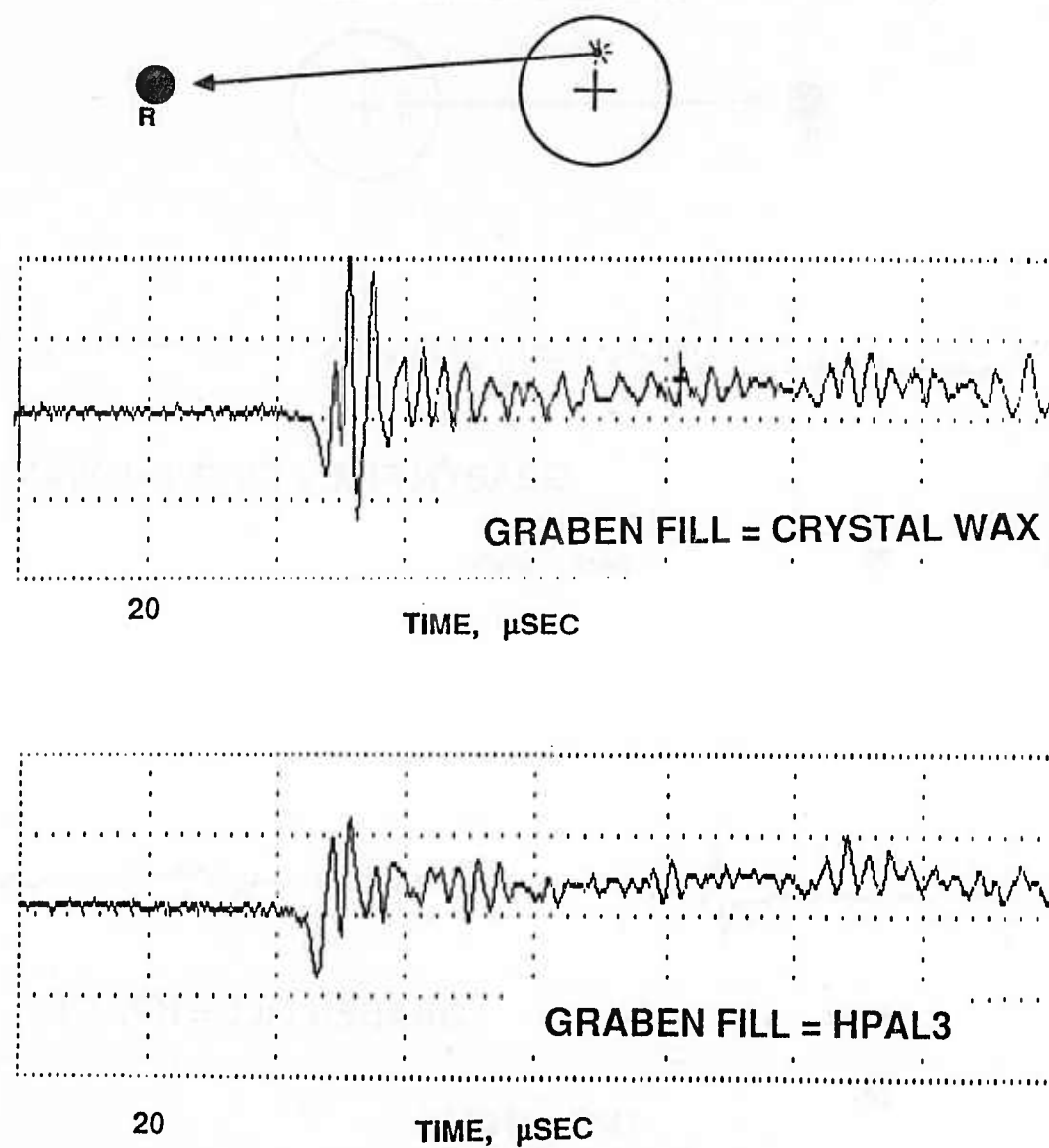


Fig. 12 Signals for one of the off-center positions in Fig. 9, for a graben filled with Crystalbond 504 ("crystal wax") and a graben filled with HPAL3. Distance from graben center to receiver is 200 mm. Vertical scale is 100 mV per division.

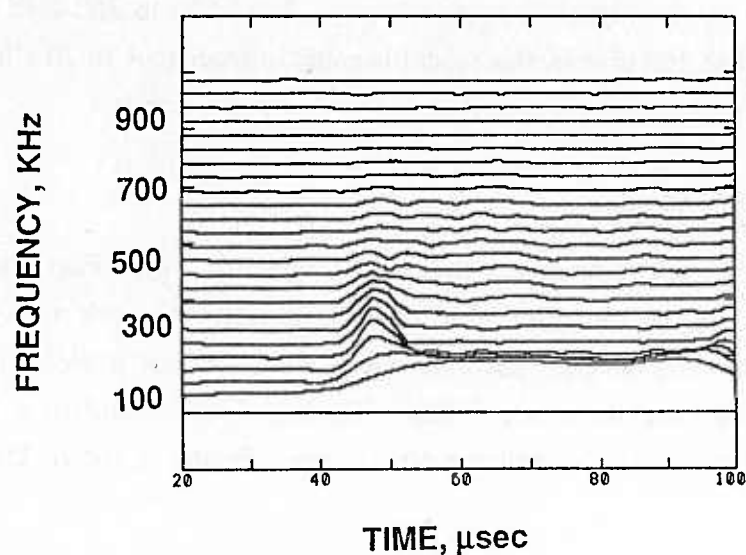


Fig. 13 "Voiceprint" of the halfspace signal shown in Fig. 6. Each trace in the voiceprint represents the signal filtered by a bandpass filter with a bandwidth of 200 kHz. The increment in center frequency of the filter is 40 kHz as we move from the bottom trace upwards. Thus, this is a time-frequency diagram.

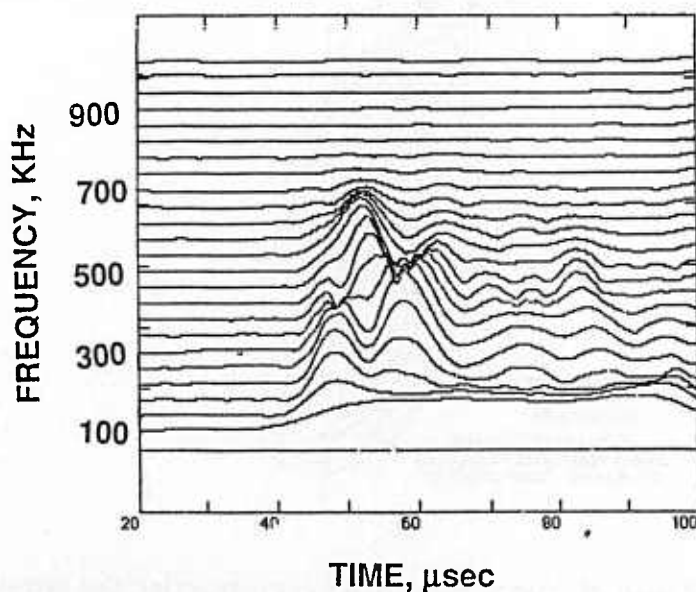


Fig. 14 Voiceprint similar to that in Fig. 13, but this time for the signal of Fig. 7, the signal from the source actuated in the graben center when the graben is filled with Crystalbond 504 ("crystal wax").

Examination of the voiceprints shows that although the low-frequency levels are quite similar between the two cases, the case with the source in the graben has considerably more energy in the higher frequency range, from 500 to 700 kHz center frequency. Considering that 100 kHz in the model is roughly analogous to 20 s in the Earth, this 500 to 700 kHz range corresponds roughly to 3 or 4 s in the Earth.

6. The Yucca Flat Model

We have constructed an accurate scale model of Yucca Flat, based on the generalized map shown in Fig. 15, which was constructed from the work of Ferguson et al (1986). The rectangular box outlines an area where many nuclear explosions have been detonated (see, e.g., McLaughlin et al, 1986). The basin was drilled in a fine-grained gabbro halfspace and filled with Crystalbond 504 (Aremco Products, Inc.). The properties

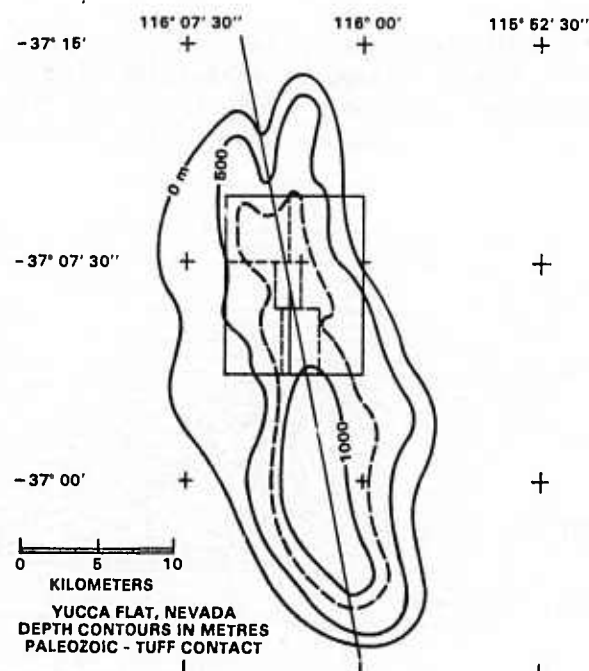


Fig. 15 Generalized map of Yucca Flat used in constructing the physical model. Map was drawn based on the work of Ferguson et al (1986). Large rectangle delineates area where many explosions were located.

of the gabbro and Crystalbond are given in Table 1. The source is excited on the surface of the basin and receivers are placed outside the basin on the halfspace, as shown in Fig. 16.

1 mm in the model represents 1 km in the Earth. If the material filling the basin in the Earth has a Rayleigh wave velocity $V_R = 1.2$ km/s, then a 20 s Rayleigh wave in the Earth has a wavelength of 24 km. Thus, the analog of the Earth's 20 s Rayleigh wave has a wavelength of 24 mm. Given that the basin fill material has a Rayleigh wave velocity of 1.01 km/s, this corresponds to a frequency of about 40 kHz. Thus, in this scale model of Yucca Flat, a 40 kHz Rayleigh wave corresponds *roughly* to 20 s Rayleigh wave in the Earth.

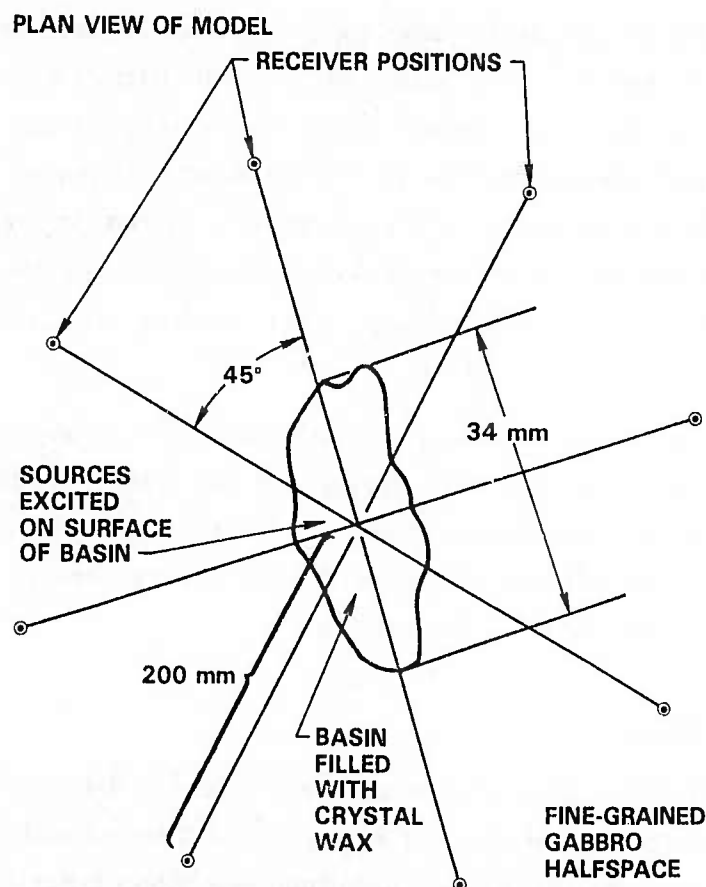


Fig. 16 Overall geometry of the model. "Crystal wax" refers to Crystalbond 504 (Aremco Prods., Inc.)

6.1 Source Excited Within the Scale Model of Yucca Flat

Figures 18-21 show the waveforms obtained when the source was excited on the surface of the basin. For comparison, Fig. 17 shows the waveforms obtained when the source was excited in the gabbro before the basin was excavated. In each case, the receiver positions are identical - 200 mm from the same reference point in the basin (the same point as the source position in Fig. 20).

The figures are meant primarily to illustrate waveform character; they should not be used as an accurate guide to arrival times. Arrival times are not completely reliable because of triggering problems in the apparatus, which we have subsequently rectified.

Examination of Figs. 18-21 yields the perhaps rather surprising result that the presence of the basin does not seem to be making much difference in the shape of the waveform. When the source is anywhere within the rectangular box representing the location of many actual explosions (Figs. 18-20), the waveforms appear almost like half-space responses. There does appear to be some complexity and ringing, but nothing like what is observed in the case of the cylindrical plug graben. The loss of symmetry in going from a cylindrical plug graben to a more realistic structure seems to have dramatically reduced the focusing effects.

The only case where relatively dramatic effects on the waveshapes is observed is shown in Fig. 21, when the source is excited over the deepest portion of the graben. Here, we see some of the same kind of complexity and ringing observed in the cylindrical plug graben. This is not difficult to rationalize, because the structure at this position locally approximates a plug or bowl-like structure.

6.2 Spectral Ratios

Figures 22 and 23 show some spectral ratios for the data in Figs. 19 and 21. At each position, the magnitude spectrum of the waveform obtained with the source excited on the basin surface is divided by the magnitude spectrum of the waveform obtained through the gabbro before the basin was excavated. If the spectral value of the denominator is too small (5% or less of its maximum value), the ratio is set to zero. Thus, we obtain some notion of the effect of the structure on the wave propagation.

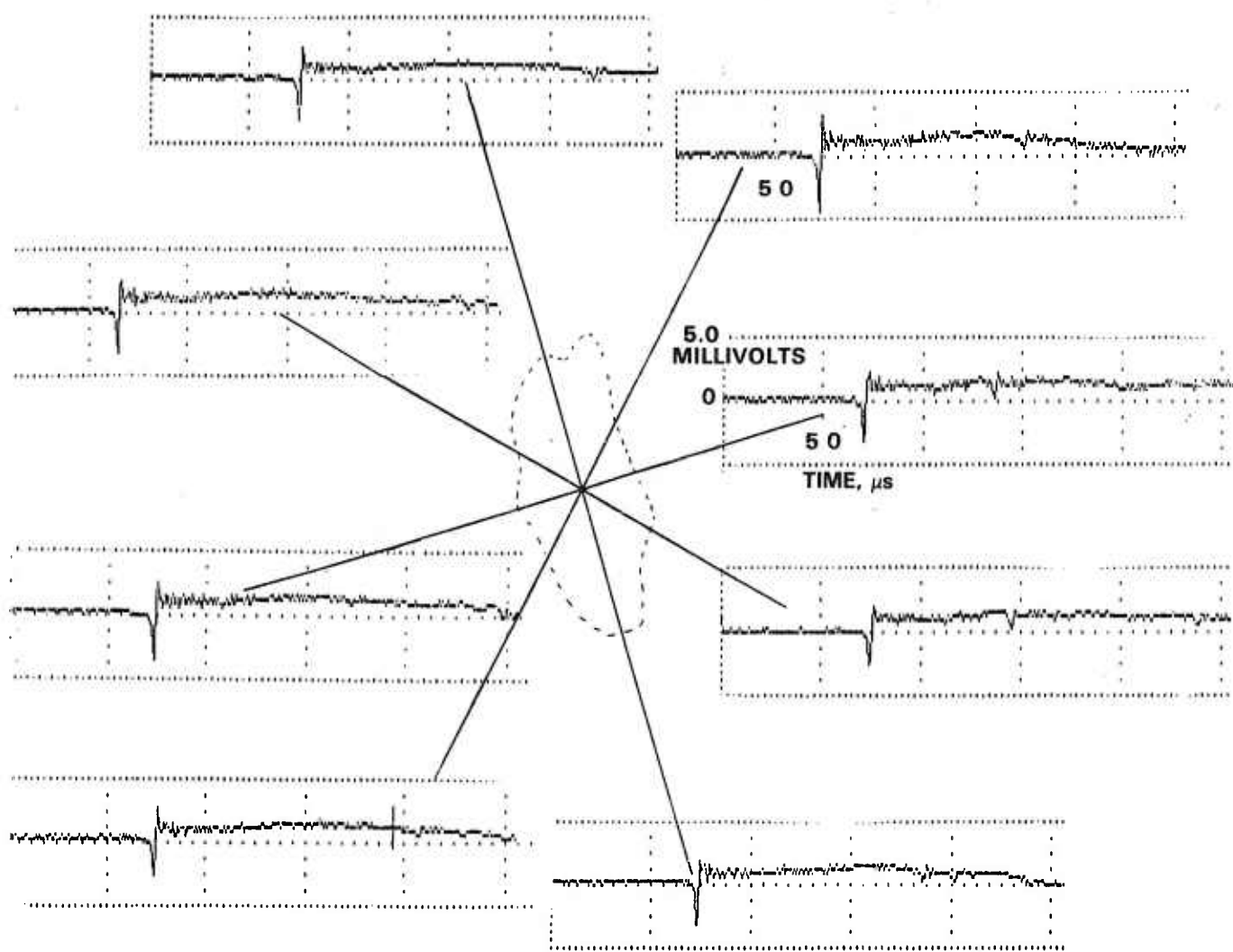


Fig. 17 Waveforms obtained on the gabbro halfspace before the Yucca Flat model was excavated. Vertical scale is 50 mV per division. Horizontal scale is 50 μ s per division.

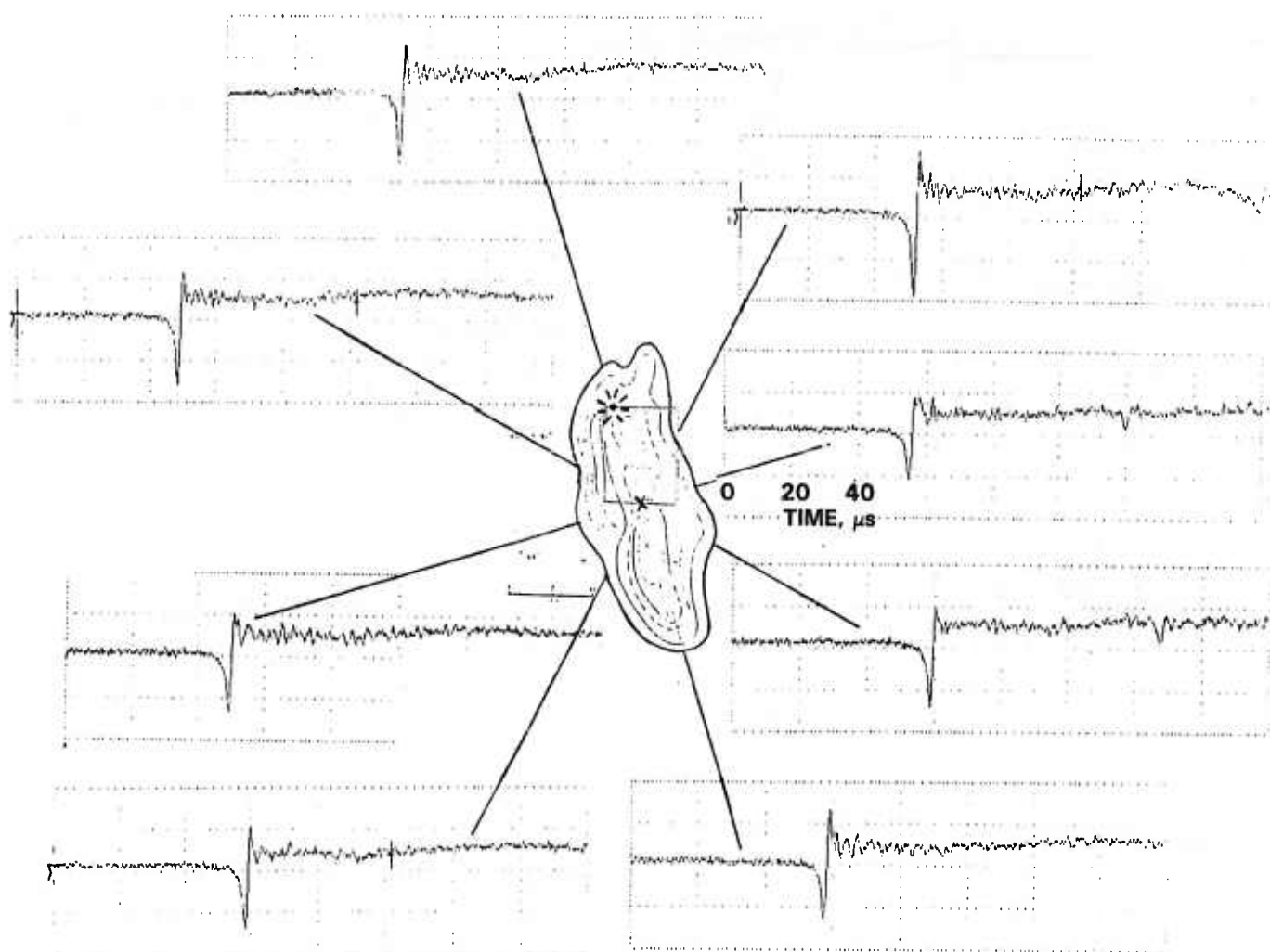


Fig. 18 Waveforms obtained from sources excited on the surface of the basin. The point "X" in each case is the reference point - all receivers are 200 mm from this reference point. The dot with the rays coming out of it indicates the position of the source in each case. In each case, the vertical scale is 20 mV per division, and the horizontal scale is 20 μ s per division.

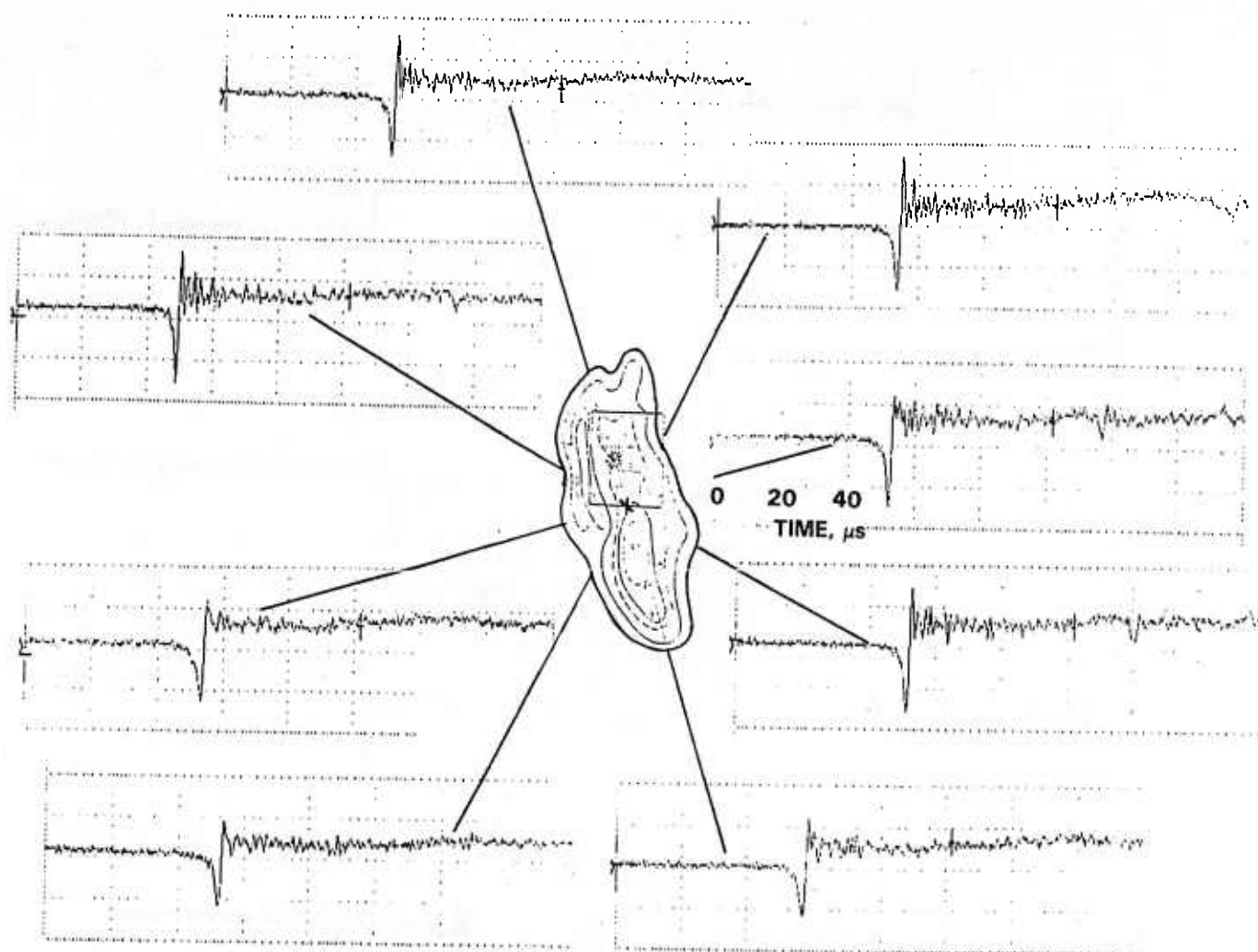


Fig. 19 Waveforms obtained from sources excited on the surface of the basin. The point "X" in each case is the reference point - all receivers are 200 mm from this reference point. The dot with the rays coming out of it indicates the position of the source in each case. In each case, the vertical scale is 20 mV per division, and the horizontal scale is 20 μs per division.

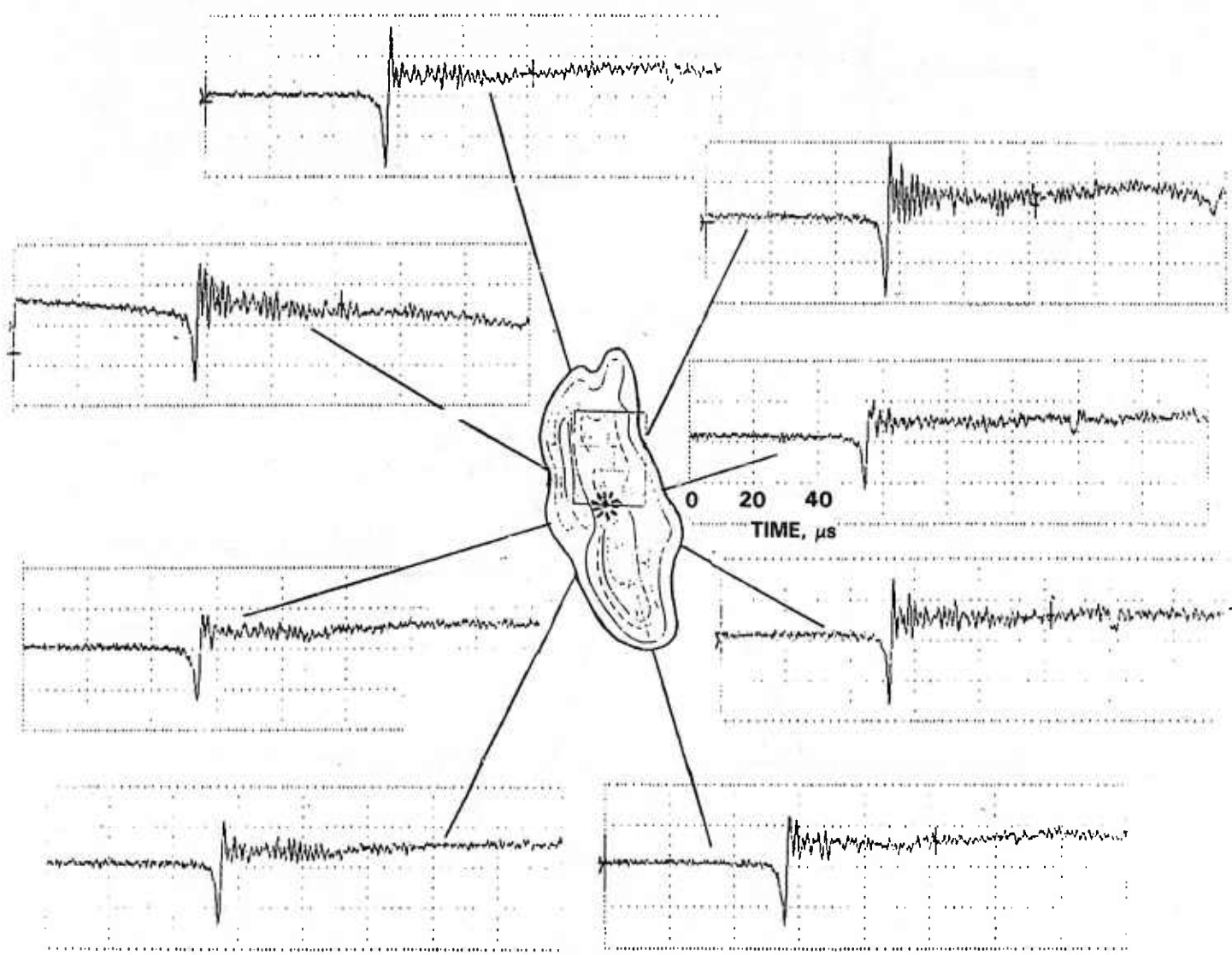


Fig. 20 Waveforms obtained from sources excited on the surface of the basin. The point "X" in each case is the reference point - all receivers are 200 mm from this reference point. The dot with the rays coming out of it indicates the position of the source in each case. In each case, the vertical scale is 20 mV per division, and the horizontal scale is 20 μs per division.

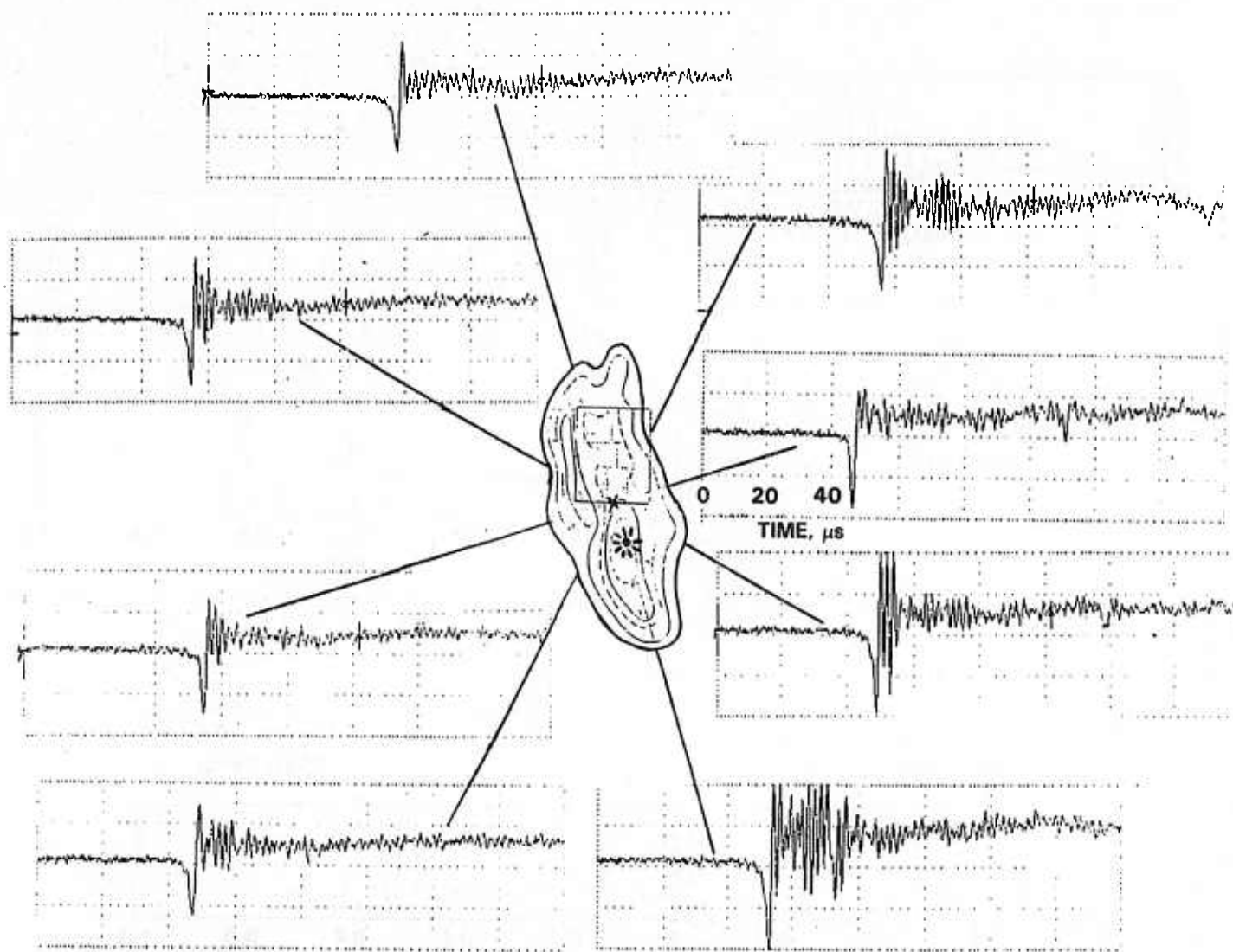


Fig. 21 Waveforms obtained from sources excited on the surface of the basin. The point "X" in each case is the reference point - all receivers are 200 mm from this reference point. The dot with the rays coming out of it indicates the position of the source in each case. In each case, the vertical scale is 20 mV per division, and the horizontal scale is 20 μs per division.

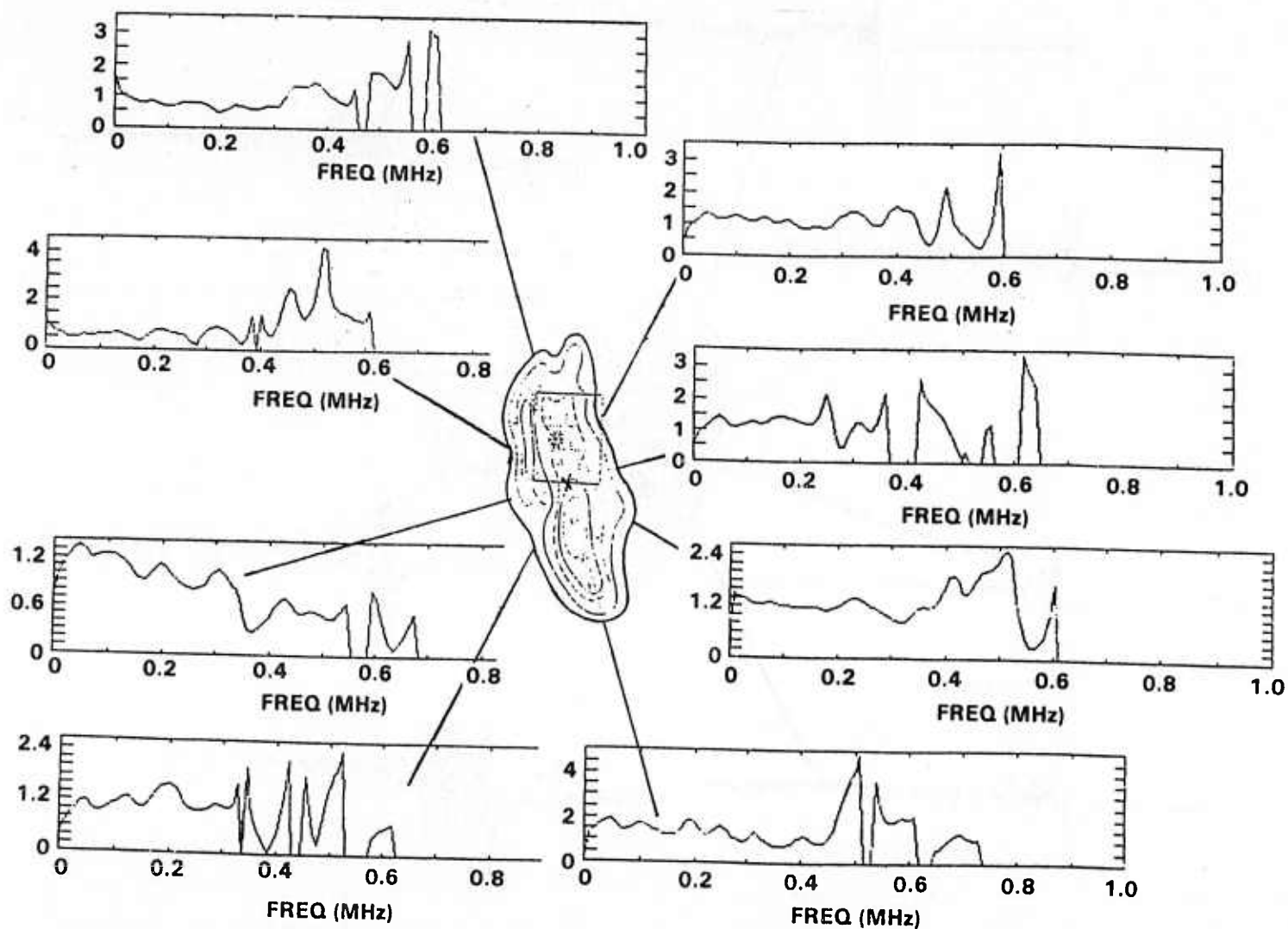


Fig. 22 Spectral ratios for the data in Fig. 19. At each position, the magnitude spectrum of the waveform obtained with the source excited on the basin surface is divided by the magnitude spectrum of the waveform obtained through the gabbro before the basin was excavated.

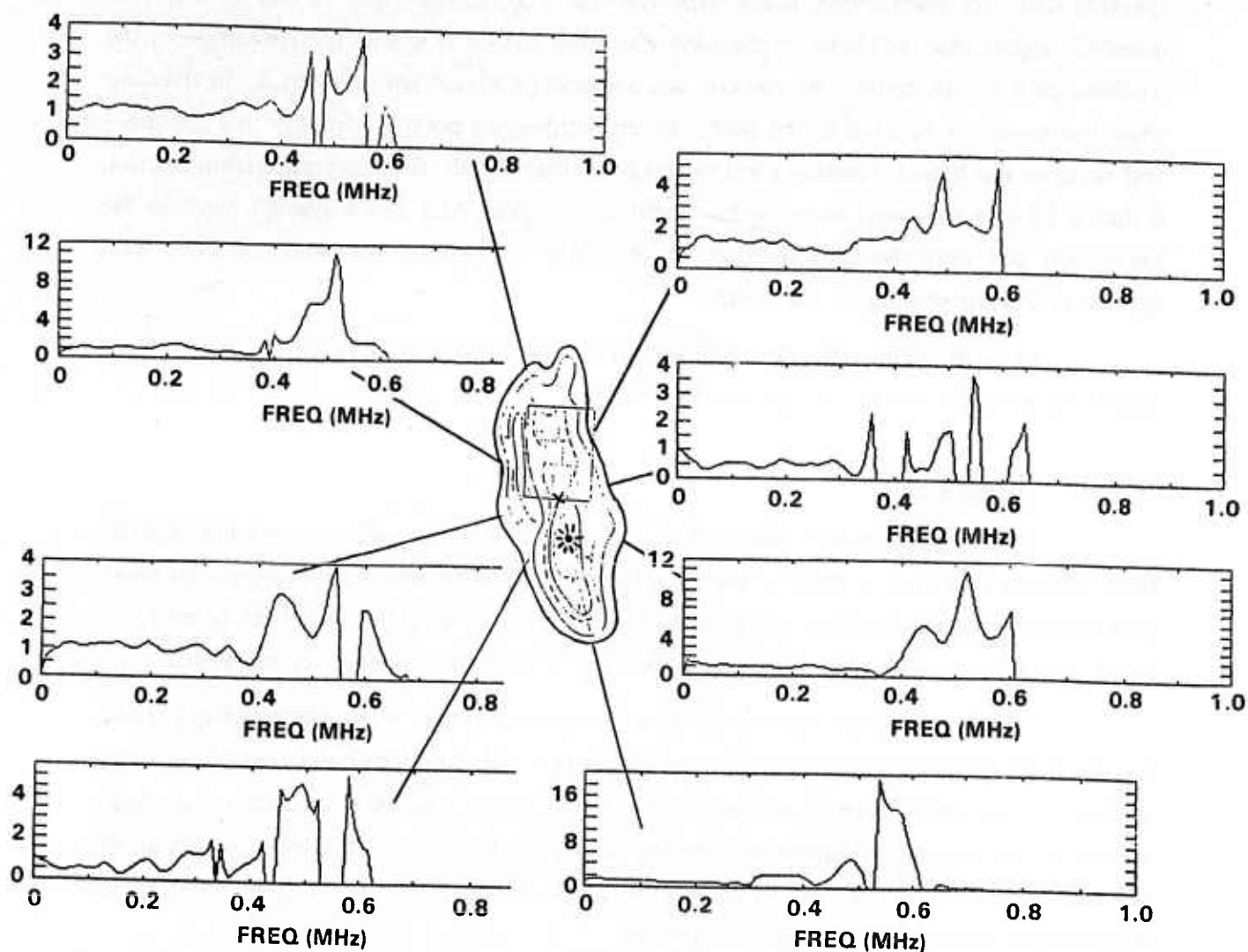


Fig. 23 Spectral ratios for the data in Fig. 21. At each position, the magnitude spectrum of the waveform obtained with the source excited on the basin surface is divided by the magnitude spectrum of the waveform obtained through the gabbro before the basin was excavated.

The spectral ratios indicate that the structure seems to be causing some amplification of high frequencies, above 400 kHz, relative to the lower frequencies. The spectral ratio for frequencies lower than 400 kHz is generally close to 1 or 2. For frequencies higher than 400 kHz, in the case when the source is within the rectangle in the northern part of the basin, the spectral ratio reaches a maximum of about 5. In the case when the source is located to the south, above the deepest portion of the basin, the spectral ratio in the high-frequency portion can be as high as 20. Remembering from Section 6 that a 40 kHz Rayleigh wave in the model is analogous to a 20 s Rayleigh wave in the Earth, we see that the amplification is occurring for waves analogous to ones with periods of 2 s and shorter in the Earth.

In both cases, the directions of greatest amplifications of high frequencies appear to lie in the southeast and the northwest.

7. Conclusions

We have described experiments intended to clarify seismic wave propagation from sources actuated in graben-like structures. We have studied two models, an idealized cylindrical graben and a scale model (1 mm to 1 km) of the Yucca Flat basin excavated into a halfspace of fine-grained gabbro and filled with low-velocity material.

Ultrasonic waves were excited on the surface of the model basin using a breaking pencil lead as a source; this source represents a step-function point unloading of the surface. The waves have been monitored using a true displacement conical transducer placed on the gabbro halfspace outside the basin, 200 mm away. Rayleigh waves of 40-120 kHz in the model correspond roughly to Rayleigh waves of period 20 s in the Earth, depending on the model and the fill material.

First, we made measurements setting the source off on the halfspace (made of gabbro, with $V_p = 6.2$ km/s), and within a cylindrical "graben" of 13 mm diameter and 2 mm depth. The graben was filled with either Crystalbond 504 ($V_p = 2.407$) or HPAL3 ($V_p = 3.287$). The presence of a source region with significantly slower velocities than the surrounding region appears to lead to a more complex signal, with more "ringing" than would be apparent if there were no such source region. The presence of such a source region appears to result in a relative amplification of the high-frequency part of the signal. The frequencies analogous to 3-10 s in the Earth appear to be amplified rela-

tive to lower frequencies. When the source is set off in the graben in an off-center position, a radiation pattern is established, with amplitude varying by a factor of 2 or more. Material effects appear to be accentuated when the source is excited off-center.

In the case of the more realistic scale model of Yucca Flat, the presence of the basin was also found to have an effect on the waveforms obtained. Some ringing and complexity are introduced into the waveforms, compared with waveforms obtained on a halfspace without a basin present. However, the effects are less dramatic than those observed when sources are excited on a model basin which is perfectly cylindrical in geometry. The most complex waveforms are obtained when the source is excited over the deepest portion of the graben. Here, we see some of the same kind of complexity and ringing observed in the cylindrical plug graben. When the source is anywhere within a rectangular box, representing the location of many actual explosions (Figs. 18-20), the effects are much less pronounced; the waveforms appear almost like halfspace responses.

The presence of the basin can cause some amplification of higher frequencies. Frequencies higher than about 400 kHz, which correspond roughly to periods of 2 s or shorter in the Earth, appear to be amplified relative to lower frequencies. This effect is most pronounced when the source is in the southern portion of the basin, as compared to the case when it is in the northern portion. The data also suggest an additional enhancement of this effect for wave propagation directions to the northwest or southeast.

In the real Earth, Yucca Flat is not embedded in a homogeneous halfspace, but in a more complex multilayered structure, and this may have a significant effect on seismic waveforms. It would be beneficial to conduct experiments involving grabens embedded in multilayered structures.

Clearly, a breaking pencil lead is a different source from a nuclear explosion, and although it is well characterized and useful in experiments such as these, it is not an exact model of a bomb. Thus, some caution should be exercised in the interpretation of these results.

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